

Learning dynamics of spring-mass models with physics-informed graph neural networks

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We propose a physics-informed message-passing graph neural network (GNN) for learning the dynamics of spring-mass systems. The proposed method embeds the underlying physics directly into the message-passing scheme of the GNN. We compare the new scheme with conventional message passing and demonstrate the generalization capability of the method. Additionally, we infer the learned parameters of the edges and show that these parameters serve as explainable metrics for the learned physics. The numerical results indicate that the proposed method accurately learns the physics of the spring-mass systems.

Keywords: Graph networks, Physics simulations, coupled-spring mass systems.

1. Introduction

Previous research has shown that GNNs can efficiently predict pairwise interactive dynamics between discrete masses when represented as graphs Battaglia and et. al. [2018] Pfaff and et al. [2021]. This is due to the inductive bias of GNNs, which models pairwise interactions as messages passed between the edges of the graph. Recently, a new type of GNNs called E(n) - Equivariant Graph Neural Networks - have been proposed Satorras and et al. [2021]. The goal of these networks is to enhance the inductive bias by preserving the rotation and translation symmetries in Euclidean space. These networks operate on Euclidean graphs where nodes and edges are assigned positions and relative distance vectors, respectively, in addition to their corresponding features. Building upon the aforementioned research, our study presents a novel message-passing scheme that integrates physical inductive biases to simulate the dynamics of spring-mass systems. Similar to the prior work, our method operates on graphs defined in Euclidean space.

2. Physics-informed GNN

We propose a GNN framework for spring-mass systems, where learned edge messages operate on embeddings of connected nodes.

2.1. Physics-informed message passing

Our message passing is based on the following ingredients. **(1)** Edge message m_{ij} is constructed as $m_{ij} = \phi_e(\|\vec{x}_{ij}\|^2)$, where ϕ_e is a learned function and $\vec{x}_{ij} = \vec{x}_i - \vec{x}_j$ is the edge vector. Our formulation results in $m_{ij} = m_{ji}$. We consider m_{ij} to be a hidden representation of the magnitude of the spring force. **(2)** At each node incoming edge vectors are weighted and summed to estimate a latent internal force on nodes, where weights are determined by ϕ_w with m_{ij} as inputs i.e. $\vec{f}_i^l = \sum_{in} \vec{x}_{ij} \phi_w(m_{ij})$. This formulation ensures that the force from node i to j is equal and opposite to that from j to i . Then, two functions $\phi_{n \rightarrow s}$ and $\phi_{n \rightarrow v}$ transform the node embeddings to a scalar α and a vector \vec{f}_{ext}^l , respectively. Next, the latent node acceleration is calculated as $\ddot{\vec{x}}^l = \alpha \vec{f}_i^l + \vec{f}_{ext}^l$.

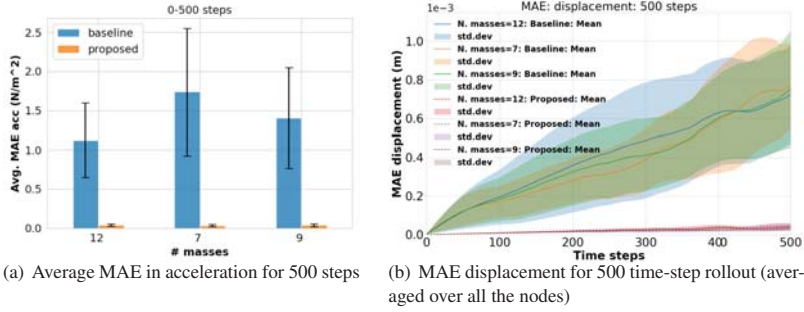


Fig. 1.: Comparison with baseline GNN

(3) The updated edge vector \vec{x}_{ij}^{upd} is calculated by adding the parallel and orthogonal projections of \vec{x}^d onto \vec{x}_{ij} . Next, the edge message is updated as $m_{ij}^{upd} = \phi_{upd}(\|\vec{x}_{ij}^{upd}\|, m_{ij})$. (4) Finally all the latent node accelerations are aggregated after n message passing steps to predict the acceleration on the node.

2.2. Inference of edge parameters

Edge parameters are inferred by calculating the eigenvalues of the Jacobian matrix of node acceleration \ddot{x}_i with respect to the edge vector \vec{x}_{ij} . For a spring with stiffness K_{ij} and rest length ℓ_{0ij} , the eigenvalues are given by:

$$eig_1 = \frac{K_{ij}(\|\vec{x}_{ij}\| - \ell_{0ij})}{\|\vec{x}_{ij}\|} \quad \text{and} \quad eig_2 = K_{ij}.$$

3. Results

To evaluate our method, we predict the trajectory roll-out of various configurations of 1kg masses connected by springs with 50N/m stiffness and 4m rest length. The GNN model proposed in Pfaff and et al. [2021] is used as baseline. Our model is trained on simulated 1000 time-step noisy trajectories of configurations with 4, 5, 6 and 8 masses. The performance of trained model was evaluated on the configurations of 7, 9 and 12 masses. Notably, the latter two configurations represent extrapolation scenarios beyond the training data. Our model achieves excellent performance in predicting the dynamics of the spring-mass systems, as well as inferring the underlying physics. The proposed method generates stable roll-outs, as shown in Figure 1 and Table 1. Table 2 shows

the inferred K and ℓ_0 parameters for each edge in a configuration of 9 masses at the 100th and 250th steps.

Table 1.: Avg MAE acceleration 0-500

	mean \pm std.	baseline	proposed
m:12 ℓ_0:4 K:50		1.286 \pm 0.207	0.125 \pm 0.096
m:3 ℓ_0:4 K:50		1.193 \pm 0.637	0.057 \pm 0.037
m:7 ℓ_0:4 K:50		1.445 \pm 0.272	0.146 \pm 0.124
m:9 ℓ_0:4 K:50		1.590 \pm 0.222	0.137 \pm 0.111

Table 2.: Inferred spring stiffness and rest length for 9-mass configuration

Edge	Mean \pm Std.dev.			
	step = 100		step = 250	
	K	ℓ_0	K	ℓ_0
0 \rightarrow 1	48.64 \pm 0.74	4 \pm 0	50.32 \pm 0.60	4.4 \pm 0.01
2 \rightarrow 1	48.52 \pm 0.62	4 \pm 0.01	49.90 \pm 0.54	4.2 \pm 0
1 \rightarrow 2	50.14 \pm 2.35	4 \pm 0	49.81 \pm 0.40	4.1 \pm 0
3 \rightarrow 2	48.46 \pm 0.69	4 \pm 0	48.76 \pm 0.98	4 \pm 0.01
2 \rightarrow 3	50.22 \pm 2.41	4 \pm 0	48.71 \pm 0.96	4 \pm 0.01
4 \rightarrow 3	50.16 \pm 2.31	4 \pm 0	49.72 \pm 0.33	4.1 \pm 0.01
3 \rightarrow 4	48.4 \pm 0.64	4 \pm 0.01	49.84 \pm 0.42	4.2 \pm 0
5 \rightarrow 4	48.41 \pm 0.67	4 \pm 0.01	49.84 \pm 0.44	4.2 \pm 0
4 \rightarrow 5	50.13 \pm 2.35	4 \pm 0	49.7 \pm 0.41	4.1 \pm 0
6 \rightarrow 5	48.46 \pm 0.69	4 \pm 0	48.69 \pm 0.96	4 \pm 0.01
5 \rightarrow 6	50.17 \pm 2.32	4 \pm 0	48.78 \pm 0.98	4 \pm 0.01
7 \rightarrow 6	50.21 \pm 2.40	4 \pm 0	49.82 \pm 0.34	4.1 \pm 0.01
6 \rightarrow 7	48.52 \pm 0.60	4 \pm 0.01	49.9 \pm 0.53	4.2 \pm 0
8 \rightarrow 7	48.57 \pm 0.71	4 \pm 0.01	50.26 \pm 0.50	4.4 \pm 0.01

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